

The Crystal Structure of a Compound with an Uncommon Aza-purine Ring System: 5,7-Bis(dimethylamino)-2-(methylthio)-*s*-triazolo-[1,5-*a*]-*s*-triazine

BY RICHARD D. GILARDI

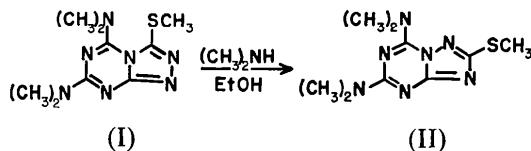
Laboratory for the Structure of Matter, Naval Research Laboratory, Washington, D.C. 20375, U.S.A.

(Received 27 April 1973; accepted 12 May 1973)

The crystal structure of 5,7-bis(dimethylamino)-2-(methylthio)-*s*-triazolo-[1,5-*a*]-*s*-triazine, C₉H₁₅N₇S, has been determined by X-ray diffraction. 1358 reflections were measured on an automatic diffractometer with Mo *K*α radiation. The space group of the crystal is *Cc*, with *a* = 9.370 (0.0024), *b* = 17.989 (0.0040), *c* = 7.842 (0.0014) Å, β = 109.84 (2)°, and 4 molecules per unit cell. The ring system of the molecule, which has not been previously characterized by diffraction methods, is planar to within 0.01 Å and resembles the purine ring system in its bond lengths and angles. The dimethylamino and methylthio ring substituents are twisted out of the plane by less than 10°. Bond lengths indicate that the dimethylamino nitrogen atoms are conjugated to the semi-aromatic ring system. The almost planar molecules assemble in approximately coplanar sheets which are separated by only 3.36 Å. The structure was solved by the symbolic addition procedure and refined to an *R* of 0.041 (calculated with all data).

Introduction

Many highly active insect chemosterilants are dimethylamino-substituted derivatives of *s*-triazine or similar heterocyclic ring systems (for references, see DeMilo, Oliver & Gilardi, 1973). Recently, during synthetic work aimed at producing additional chemosterilants, an unexpected rearrangement of the *s*-triazolo-[4,3-*a*]-*s*-triazine ring system was encountered (DeMilo, Oliver & Gilardi, 1973). A synthetic intermediate (I) containing this ring system was treated with excess anhydrous dimethylamine in absolute ethanol. The expected dimethylamino replacement at the thio-methyl site did not occur; analysis of the product indicated, rather, that an isomerization had occurred. The most probable explanation for this isomerization seemed to be a rearrangement in the ring system, yielding a derivative of *s*-triazolo-[1,5-*a*]-*s*-triazine (II).



The proposed structure (II) of the product, 5,7-bis(dimethylamino)-*s*-(methylthio)-*s*-triazolo-[1,5-*a*]-*s*-triazine, is corroborated by the X-ray diffraction analysis reported in this article. The ring system is very similar to that of the purine bases, and this compound could, in fact, be named as a 5-azapurine derivative, with the conventional purine numbering scheme. However, in this paper, the compound is named and numbered as a triazolo-triazine so as to conform with existing references in *Chemical Abstracts*. Structural parameters for the heterocyclic ring system in this molecule have not been previously reported.

Experimental

Crystalline samples of (II) were supplied by Dr A. B. DeMilo of the USDA Agricultural Research Center (Beltsville, Md.). The compound was recrystallized by evaporation from an acetone-cyclohexane mixture. Intensities were collected from a single crystal at room temperature with Mo *K*α radiation (0.71069 Å) on a fully automated computer-controlled 4-circle diffractometer with a θ-2θ scan and a 1.4° + 2θ(α₂) - 2θ(α₁) scan over 2θ. The crystallographic constants are reported in Table 1.

Table 1. *Physical data*

Molecular formula	C ₉ H ₁₅ N ₇ S
M.W.	253.33
Habit	Colorless, monoclinic prisms
Size	0.065 × 0.30 × 0.50 mm
Space group	<i>Cc</i>
<i>a</i>	9.370 (0.0024) Å
<i>b</i>	17.989 (0.0040)
<i>c</i>	7.842 (0.0014)
β	109.84 (0.02)°
Volume	1243 Å ³
Density (calc.)	1.353 g cm ⁻³
Density (obs.)	1.35
Molecules per unit cell	4
Radiation	Mo <i>K</i> α (λ = 0.71069 Å)
Number of independent reflections measured	1358
Maximum sin θ/λ	0.639
Linear absorption coefficient	2.48 cm ⁻¹

The intensities were corrected for Lorentz and polarization factors and placed on an absolute scale by means of a *K*-curve. Structure-factor magnitudes, |*F*_h|, and also normalized structure-factor magnitudes, |*E*_h|

were derived. The variance to be associated with each $|F_h|$ was evaluated in the first stages of data reduction, as described below.

The first step in the data reduction procedure consists of obtaining an uncorrected initial intensity value for each reflection from the counting measurements:

$$I = [P - t_f(B_1 + B_2)]$$

where P , B_1 , and B_2 refer to the peak (integrated scan) and two background measurements expressed in actual counts, and t_f is a time factor necessary to relate the fixed background counting time to the variable integrated peak scan time. The random error in the time factor is very small and is considered negligible. The variance of each counting measurement, whether peak or background, is taken to be:

$$\sigma^2(N) = N + (cN)^2$$

in which the first term is the counting error assuming Poisson statistics, while the second accounts for random intensity-dependent errors due to source or system instability. The parameter c is thus a characteristic of each particular diffraction experiment, and may be estimated from a consideration of the sample variances of the monitored standard reflections. It was calculated to be 0.01 for this particular data collection. The overall expression for the variance in the intensity is then:

$$\sigma_I^2 = P + c^2 P^2 + (t_f)^2 [B_1 + B_2 + c^2 B_1^2 + c^2 B_2^2].$$

The structure factor modulus is obtained from the intensity by:

$$|F_h| = \left(\frac{Q}{Lp} \cdot I \right)^{1/2}, \quad I > 0$$

$$|F_h| = 0, \quad I \leq 0$$

where Lp is the Lorentz-polarization factor and Q is an attenuation factor which must be used if the source was attenuated by a metal-foil filter for the measurement of intense diffraction peaks. The determination of the attenuator factors by repetitive scans of chosen peaks leads to a distribution of values with appreciable sample variances. Therefore, estimated values for σ_Q^2 were included as a source of error in the structure factor. The resultant expression for the estimate of the structure-factor variance is:

$$\sigma_{|F_h|}^2 = \frac{Q}{4Lp} \left[\left(\frac{\sigma_I}{I} \right)^2 + \left(\frac{\sigma_Q}{Q} \right)^2 \right].$$

In our experience, the two error terms in this equation are often similar in magnitude for very strong reflections. Of course, for the majority of the data, no attenuators are used, and the second term vanishes. If a reflection is very weak, the above expression may be unsuitable. If I approaches zero or is slightly negative, the first term diverges or goes negative. In practice, if $I < 5\sigma_I$, the following numerical approximation gives reasonable values for $\sigma_{|F_h|}$ and has been applied:

$$3\sigma_{|F_h|} = \left[\frac{1}{Lp} (I + 3\sigma_I) \right]^{1/2} - [I/Lp]^{1/2}, \quad I \geq 0$$

$$3\sigma_{|F_h|} = \left[\frac{1}{Lp} (I + 3\sigma_I) \right]^{1/2}, \quad I < 0$$

where the first quantity on the right hand side of the equation corresponds to a (3σ) upper limit on $|F|$, and the second quantity is the expected value for $|F|$. If the intensity, I , is less than $-3\sigma_I$ there is good reason to doubt the credibility of at least one of the counting measurements, so the data reduction program rejects such reflections. The numerical approximation for $\sigma_{|F_h|}$ gives results virtually indistinguishable from the analytic form if $I > 5\sigma_I$ and could be used for all data. A similar expression was used by Sletten, Sletten & Jensen (1969) to estimate $\sigma_{|F_h|}$ throughout the entire intensity range.

Structure determination

Precession and Weissenberg photographs of this crystal indicated the space group to be either Cc (acentric) or $C2/c$ (centric). An examination of the statistical distribution of intensities did not resolve this ambiguity. For example, the value of $\langle |E| \rangle$ for the observed reflections was 0.82 *versus* theoretical values of 0.80 for a centric distribution and 0.89 for an acentric one. This would seem to indicate a centric crystal; however, the observed value of $\langle |E|^2 - 1 \rangle$ was 0.84, and the theoretical values for this statistical measure are 0.97 (centric) and 0.74 (acentric). Also, since the density indicated $Z=4$, the asymmetric unit would be only half a molecule if the space group were $C2/c$, and the molecule would possess either an inversion center or a twofold axis. This did not appear to be likely from the chemical evidence.

Thus, the phase determination was begun in space group Cc by the symbolic addition procedure (Karle & Karle, 1966). However, choosing Cc as the most likely space group does not preclude the solution of a $C2/c$ structure when this method is used for phase determination. The origin can be set in a way that is compatible with either space group, and the initial symbolic phasing procedure is the same in either case.

In space group Cc , the $0k0$ reflections have real phases and are structure seminvariants. Such reflections are therefore useful for symbolic assignments but may not be used to assign the origin. Two reflections of any other type may be used to fix the origin, but a unique assignment is assured only if

$$\begin{vmatrix} h_1 l_1 \\ h_2 l_2 \end{vmatrix} = \pm 1.$$

This condition for origin setting was originally expressed by Karle & Hauptman (1961), but in terms of indices referred to a primitive cell. Above, all references are to indices of the conventional centered cell. In this determination, the 355 and the 243 reflections were assigned zero phases to fix the origin. This pair of reflec-

tions, which obeys the above criterion, necessarily obeys the less stringent requirements for fixing the origin of a *C2/c* cell (Hauptman & Karle, 1959).

Maps were generated corresponding to the four most likely sets of phases. One map, corresponding to a set which contained both real and imaginary phases (thus, a *Cc* structure), contained a structural fragment which looked promising but not entirely satisfactory. A possible *s*-triazine ring was apparent, with the N atoms slightly heavier than the C atoms, but the S atom position was not apparent. However, the *s*-triazine ring and all atoms directly connected to it were used as a partial structure in a tangent-formula phasing procedure (Karle, 1968).

Each cycle of this procedure consists of phasing the reflections with strongest $|E|$ values from a structural fragment, refining those phases and determining additional phases with the tangent formula, and generating an improved *E* map from these phases. After three cycles,

the full ring system and the amino N atoms were apparent, and the S atom appeared as the heaviest atom. All methyl carbon atoms were found in a conventional difference map. A re-examination of the original *E* map revealed that the peak for the S atom was indeed present, but was seventh in intensity, and could not be connected to the partial structure because C(2), the adjacent ring atom, was missing in that first map.

An examination of the coordinates of the refined structure revealed a pseudo-center of the molecule, which could explain some of the ambiguities encountered in the determination of this structure. A shift of $\Delta x = -0.517$, $\Delta z = -0.895$ (arbitrary shifts in *x* and *z* preserve the crystallographic symmetry in *Cc*) moves the center of the *s*-triazine ring to the origin. The positions of the atoms of the *s*-triazine ring, as well as N(1) and N(13), are now quite accurately related by an inversion center at the origin. [C(2) and C(14) also constitute an approximate inversion pair]. If the other

Table 2. Observed and calculated structure factors for 5,7-bis(dimethylamino)-2-(methylthio)-*s*-triazolo-[1,5-*a*]-*s*-triazine (II)

The columns are the index *hkl*, $10|F_o|$, and $10|F_c|$.

<i>hkl</i>	$10 F_o $	$10 F_c $	<i>hkl</i>	$10 F_o $	$10 F_c $	<i>hkl</i>	$10 F_o $	$10 F_c $	<i>hkl</i>	$10 F_o $	$10 F_c $
0 0 0	0	0	0 0 0	0	0	0 0 0	0	0	0 0 0	0	0
2 000	114	114	0 20 0	111	111	0 0 20	111	111	0 0 20	111	111
4 040	237	237	0 40 0	222	222	0 0 40	222	222	0 0 40	222	222
6 130	54	54	0 60 0	333	333	0 0 60	333	333	0 0 60	333	333
8 270	120	120	0 80 0	444	444	0 0 80	444	444	0 0 80	444	444
10 490	267	267	0 100 0	555	555	0 0 100	555	555	0 0 100	555	555
12 890	594	594	0 120 0	666	666	0 0 120	666	666	0 0 120	666	666
14 1660	1311	1311	0 140 0	777	777	0 0 140	777	777	0 0 140	777	777
16 2990	2934	2934	0 160 0	888	888	0 0 160	888	888	0 0 160	888	888
18 5380	6486	6486	0 180 0	999	999	0 0 180	999	999	0 0 180	999	999
20 9820	14550	14550	0 200 0	1100	1100	0 0 200	1100	1100	0 0 200	1100	1100
22 17740	32112	32112	0 220 0	1211	1211	0 0 220	1211	1211	0 0 220	1211	1211
24 32540	71826	71826	0 240 0	1322	1322	0 0 240	1322	1322	0 0 240	1322	1322
26 59820	161106	161106	0 260 0	1433	1433	0 0 260	1433	1433	0 0 260	1433	1433
28 109440	358224	358224	0 280 0	1544	1544	0 0 280	1544	1544	0 0 280	1544	1544
30 201840	792462	792462	0 300 0	1655	1655	0 0 300	1655	1655	0 0 300	1655	1655
32 375120	1731422	1731422	0 320 0	1766	1766	0 0 320	1766	1766	0 0 320	1766	1766
34 682380	3822666	3822666	0 340 0	1877	1877	0 0 340	1877	1877	0 0 340	1877	1877
36 1248660	8418186	8418186	0 360 0	1988	1988	0 0 360	1988	1988	0 0 360	1988	1988
38 2286120	18520824	18520824	0 380 0	2099	2099	0 0 380	2099	2099	0 0 380	2099	2099
40 4201800	40546662	40546662	0 400 0	2210	2210	0 0 400	2210	2210	0 0 400	2210	2210
42 7743600	89203322	89203322	0 420 0	2321	2321	0 0 420	2321	2321	0 0 420	2321	2321
44 14118000	195206644	195206644	0 440 0	2432	2432	0 0 440	2432	2432	0 0 440	2432	2432
46 26111400	428222288	428222288	0 460 0	2543	2543	0 0 460	2543	2543	0 0 460	2543	2543
48 47822400	942444576	942444576	0 480 0	2654	2654	0 0 480	2654	2654	0 0 480	2654	2654
50 88233600	2062222222	2062222222	0 500 0	2765	2765	0 0 500	2765	2765	0 0 500	2765	2765
52 160464000	4522222222	4522222222	0 520 0	2876	2876	0 0 520	2876	2876	0 0 520	2876	2876
54 296832000	9922222222	9922222222	0 540 0	2987	2987	0 0 540	2987	2987	0 0 540	2987	2987
56 544320000	2162222222	2162222222	0 560 0	3098	3098	0 0 560	3098	3098	0 0 560	3098	3098
58 1000640000	4662222222	4662222222	0 580 0	3209	3209	0 0 580	3209	3209	0 0 580	3209	3209
60 1860960000	1016222222	1016222222	0 600 0	3320	3320	0 0 600	3320	3320	0 0 600	3320	3320
62 3441280000	2216222222	2216222222	0 620 0	3431	3431	0 0 620	3431	3431	0 0 620	3431	3431
64 6342560000	4816222222	4816222222	0 640 0	3542	3542	0 0 640	3542	3542	0 0 640	3542	3542
66 11785120000	1041622222	1041622222	0 660 0	3653	3653	0 0 660	3653	3653	0 0 660	3653	3653
68 21970240000	2261622222	2261622222	0 680 0	3764	3764	0 0 680	3764	3764	0 0 680	3764	3764
70 41112480000	4921622222	4921622222	0 700 0	3875	3875	0 0 700	3875	3875	0 0 700	3875	3875
72 76424960000	1062162222	1062162222	0 720 0	3986	3986	0 0 720	3986	3986	0 0 720	3986	3986
74 141850880000	2302162222	2302162222	0 740 0	4097	4097	0 0 740	4097	4097	0 0 740	4097	4097
76 265701760000	4962162222	4962162222	0 760 0	4208	4208	0 0 760	4208	4208	0 0 760	4208	4208
78 496003520000	1076216222	1076216222	0 780 0	4319	4319	0 0 780	4319	4319	0 0 780	4319	4319
80 920007040000	2316216222	2316216222	0 800 0	4430	4430	0 0 800	4430	4430	0 0 800	4430	4430
82 1716014080000	4976216222	4976216222	0 820 0	4541	4541	0 0 820	4541	4541	0 0 820	4541	4541
84 3216028160000	1077621622	1077621622	0 840 0	4652	4652	0 0 840	4652	4652	0 0 840	4652	4652
86 5968056320000	2317621622	2317621622	0 860 0	4763	4763	0 0 860	4763	4763	0 0 860	4763	4763
88 11056112640000	4977621622	4977621622	0 880 0	4874	4874	0 0 880	4874	4874	0 0 880	4874	4874
90 20512225280000	10777621622	10777621622	0 900 0	4985	4985	0 0 900	4985	4985	0 0 900	4985	4985
92 38024450560000	23177621622	23177621622	0 920 0	5096	5096	0 0 920	5096	5096	0 0 920	5096	5096
94 70048901120000	49777621622	49777621622	0 940 0	5207	5207	0 0 940	5207	5207	0 0 940	5207	5207
96 131697802240000	107777621622	107777621622	0 960 0	5318	5318	0 0 960	5318	5318	0 0 960	5318	5318
98 246400004480000	231777621622	231777621622	0 980 0	5429	5429	0 0 980	5429	5429	0 0 980	5429	5429
100 458400008960000	497777621622	497777621622	0 1000 0	5540	5540	0 0 1000	5540	5540	0 0 1000	5540	5540
102 852800017920000	1077777621622	1077777621622	0 1020 0	5651	5651	0 0 1020	5651	5651	0 0 1020	5651	5651
104 1569600035840000	2317777621622	2317777621622	0 1040 0	5762	5762	0 0 1040	5762	5762	0 0 1040	5762	5762
106 2918400071680000	4977777621622	4977777621622	0 1060 0	5873	5873	0 0 1060	5873	5873	0 0 1060	5873	5873
108 5347200143360000	10777777621622	10777777621622	0 1080 0	5984	5984	0 0 1080	5984	5984	0 0 1080	5984	5984
110 9984000286720000	23177777621622	23177777621622	0 1100 0	6095	6095	0 0 1100	6095	6095	0 0 1100	6095	6095
112 18528000573440000	49777777621622	49777777621622	0 1120 0	6206	6206	0 0 1120	6206	6206	0 0 1120	6206	6206
114 34464001146880000	107777777621622	107777777621622	0 1140 0	6317	6317	0 0 1140	6317	6317	0 0 1140	6317	6317
116 63840002293760000	231777777621622	231777777621622	0 1160 0	6428	6428	0 0 1160	6428	6428	0 0 1160	6428	6428
118 118560004587520000	497777777621622	497777777621622	0 1180 0	6539	6539	0 0 1180	6539	6539	0 0 1180	6539	6539
120 220320009175040000	1077777777621622	1077777777621622	0 1200 0	6650	6650	0 0 1200	6650	6650	0 0 1200	6650	6650
122 412800018350080000	2317777777621622	2317777777621622	0 1220 0	6761	6761	0 0 1220	6761	6761	0 0 1220	6761	6761
124 770400036700160000	4977777777621622	4977777777621622	0 1240 0	6872	6872	0 0 1240	6872	6872	0 0 1240	6872	6872
126 1417600073400320000	10777777777621622	10777777777621622	0 1260 0	6983	6983	0 0 1260	6983	6983	0 0 1260	6983	6983
128 2649600146800640000	23177777777621622	23177777777621622	0 1280 0	7094	7094	0 0 1280	7094	7094	0 0 1280	7094	7094
130 49248002936001280000	49777777777621622	49777777777621622	0 1300 0	7205	7205	0 0 1300	7205	7205	0 0 1300	7205	7205
132 91536005872002560000	107777777777621622	107777777777621622	0 1320 0	7316	7316	0 0 1320	7316	7316	0 0 1320	7316	7316
134 1710720011744005120000	231777777777621622	231777777777621622	0 1340 0	7427	7427	0 0 1340	7427	7427	0 0 1340	7427	7427
136 32214400234880010240000	497777777777621622	497777777777621622	0 1360 0	7538	7538	0 0 1360	7538	7538	0 0 1360	7538	7538
138 5980800469760020480000	1077777777777621622	1077777777777621622	0 1380 0	7649	7649	0 0 1380	7649	7649	0 0 1380	7649	7649
140 11073600939520040960000	2317777777777621622	2317777777777621622	0 1400 0	7760	7760	0 0 1400	7760	7760	0 0 1400	7760	7760
142 208224001879040081920000	4977777777777621622	4977777777777621622	0 1420 0	7871	7871	0 0 1420	7871	7871	0 0 1420	7871	7871
144 3904480037580800163840000	10777777777777621622	10777777777777621622	0 1440 0	7982	7982	0 0 1440	7982	7982	0 0 1440	7982	7982
146 7288960075161600327680000	23177777777777621622	23177777777777621622	0 1460 0	8093	8093	0 0 1460	8093	8093	0 0 1460	8093	8093
148 13891840150323200655360000	49777777777777621622	49777777777777621622	0 1480 0	8204	8204	0 0 1480	8204	8204	0 0 1480	8204	8204
150 260236803006464001310720000	107777777777777621622	107777777777777621622	0 1500 0	8315	8315	0 0 1500	8315	8315	0 0 1500	8315	8315
152 492473606012928002621440000</											

atoms were absent, this would be a true crystallographic center, and the space group would be $C2/c$. The initial phase determination gave phases which tended to reinforce this pseudo-center, and this led to an E map which emphasized those atoms that obeyed the pseudo symmetry. A re-examination of the first E map revealed that the largest spurious peak was a pseudo-inversion equivalent of the sulfur atom. This, of course, disappeared as the truly acentric portion of the structure was added and more of the phases were refined to their true acentric values.

Refinement

Isotropic and anisotropic full-matrix least-squares refinement of the non-hydrogen atoms led to an unweighted R of 0.069 at which time a difference Fourier map revealed the 15 hydrogen atoms of the structure. The hydrogen atom positions were subsequently varied, while their anisotropic thermal parameters were included in the refinement as constrained parameters; *i.e.*, the thermal parameters of all atoms of a methyl group were set equal, and varied in unison. Spherical atomic scattering factors for S, C, and N were taken from *International Tables for X-ray Crystallography* (1962), and the spherical scattering factors for bonded hydrogen of Stewart, Davidson & Simpson (1965) were used. Anomalous scattering contributions were included for S: $\Delta f' = 0.110$, $\Delta f'' = 0.124$ (Cromer & Liberman, 1970). An isotropic secondary extinction correction was applied to the calculated structure factors, as outlined by Larson (1967). Only three of the corrections were greater than 5%; the correction factors applied to F_c for the $20\bar{2}$, 110, and 020 reflections were 0.74, 0.93, and 0.91, respectively. The function minimized was $\sum w|F_o - F_c|^2$, and the final measures of agreement were

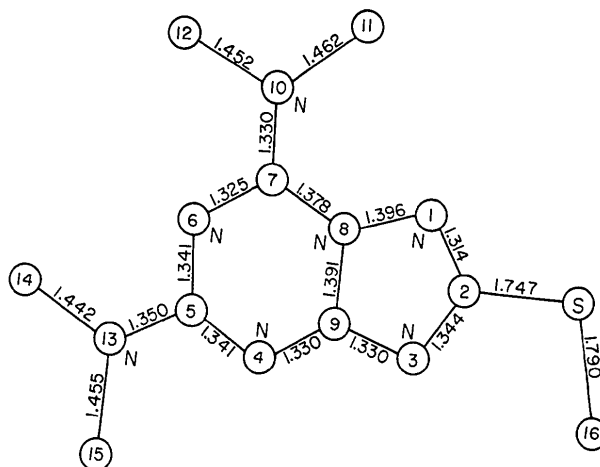


Fig. 1. Bond distances for (II). The standard deviations for bond distances indicated by the least-squares refinement are 0.005 Å for bonds to peripheral methyl carbons and 0.003–0.004 Å for all other bonds between non-hydrogen atoms.

$$R = \frac{\sum |F_o - F_c|}{\sum |F_o|} = 0.041$$

$$R_w = \frac{\sum w|F_o - F_c|^2}{\sum wF_o^2} = 0.031$$

and

$$s = \left[\frac{\sum w|F_o - F_c|^2}{N_o - N_p} \right]^{1/2} = 2.23$$

wherein N_o was 1358 and N_p , the total number of varied parameters, was 198. No reflections were considered 'unobserved' and omitted from the above summations. Final observed and calculated structure factors are given in Table 2. Final positions and thermal

Table 3. Fractional coordinates and thermal parameters with standard deviations

The thermal parameters are of the form:

$$T = \exp \left[-\frac{1}{4}(B_{11}h^2a^{*2} + B_{22}k^2b^{*2} + B_{33}l^2c^{*2} + 2B_{12}hka^*b^* + 2B_{13}hla^*c^* + 2B_{23}klb^*c^*) \right].$$

Hydrogen atoms were assigned thermal factors equal to those of the adjoining carbon atom. Standard deviations are those indicated by the full-matrix least-squares refinement.

	x	y	z	B_{11}	B_{22}	B_{33}	B_{12}	B_{13}	B_{23}
S	0.87260	0.81602 (5)	1.25440	4.76 (4)	4.46 (4)	5.36 (4)	1.27 (4)	0.87 (3)	1.25 (4)
N(1)	0.77301 (31)	0.95264 (14)	1.16039 (38)	3.73 (13)	3.77 (13)	4.03 (13)	0.57 (10)	0.69 (10)	0.42 (10)
C(2)	0.74242 (36)	0.88173 (18)	1.12764 (41)	3.72 (14)	3.80 (14)	3.76 (14)	0.70 (11)	1.28 (12)	0.47 (11)
N(3)	0.61480 (31)	0.86382 (14)	0.99160 (37)	4.11 (13)	3.36 (12)	3.98 (13)	0.54 (10)	1.03 (10)	0.26 (11)
N(4)	0.42530 (30)	0.94303 (15)	0.79491 (35)	3.75 (13)	3.51 (13)	3.74 (13)	0.29 (9)	0.69 (10)	0.38 (10)
C(5)	0.39402 (43)	1.01559 (18)	0.76481 (54)	3.53 (17)	4.15 (14)	3.83 (13)	0.67 (13)	1.51 (13)	0.69 (13)
N(6)	0.47782 (34)	1.07262 (14)	0.85595 (41)	3.68 (12)	3.27 (12)	4.14 (12)	0.47 (11)	1.27 (10)	0.53 (11)
C(7)	0.60571 (40)	1.05855 (16)	0.99063 (47)	3.94 (14)	3.46 (14)	4.17 (15)	0.22 (12)	2.01 (13)	0.08 (13)
N(8)	0.64838 (31)	0.98545 (14)	1.02974 (37)	3.57 (12)	3.22 (12)	3.66 (11)	0.31 (10)	0.98 (9)	0.03 (10)
C(9)	0.55432 (41)	0.92957 (18)	0.92965 (47)	4.37 (16)	3.39 (13)	3.62 (13)	0.40 (11)	1.33 (12)	0.12 (11)
N(10)	0.68833 (31)	1.11439 (15)	1.08509 (38)	4.13 (13)	3.10 (12)	5.62 (15)	0.17 (11)	1.53 (11)	-0.32 (11)
C(11)	0.84109 (48)	1.10938 (25)	1.21865 (68)	4.11 (18)	4.80 (21)	6.59 (22)	-0.42 (15)	0.87 (16)	-1.40 (17)
C(12)	0.62784 (57)	1.18919 (23)	1.04606 (75)	5.26 (22)	3.87 (18)	8.70 (26)	0.14 (16)	1.64 (22)	-0.75 (17)
N(13)	0.26551 (34)	1.03384 (16)	0.62901 (40)	4.06 (14)	4.48 (14)	4.35 (14)	0.42 (11)	0.63 (12)	0.44 (11)
C(14)	0.22310 (57)	1.11051 (28)	0.58715 (68)	5.48 (24)	5.73 (21)	5.84 (22)	1.55 (18)	0.86 (20)	1.55 (21)
C(15)	0.16104 (52)	0.97669 (27)	0.52791 (66)	4.66 (19)	6.75 (25)	4.62 (17)	-0.46 (20)	0.25 (16)	1.11 (18)
C(16)	0.77394 (58)	0.73198 (22)	1.16325 (74)	7.26 (26)	3.48 (16)	8.35 (28)	1.18 (18)	2.45 (24)	1.10 (20)

Table 3 (cont.)

	x	y	z
H(11A)	0.9052 (36)	1.1442 (18)	1.1680 (46)
H(11B)	0.8854 (40)	1.0660 (18)	1.2035 (49)
H(11C)	0.8476 (38)	1.1349 (19)	1.3199 (50)
H(12A)	0.6376 (39)	1.2097 (17)	1.1598 (51)
H(12B)	0.7001 (42)	1.2190 (20)	1.0396 (58)
H(12C)	0.5375 (45)	1.1906 (19)	1.0200 (58)
H(14A)	0.1590 (43)	1.1168 (22)	0.6360 (49)
H(14B)	0.2000 (42)	1.1146 (20)	0.4771 (49)
H(14C)	0.2861 (44)	1.1449 (24)	0.6121 (53)
H(15A)	0.1843 (42)	0.9312 (23)	0.5578 (50)
H(15B)	0.0732 (41)	0.9740 (19)	0.5491 (47)
H(15C)	0.1171 (39)	0.9954 (21)	0.4145 (49)
H(16A)	0.6927 (43)	0.7312 (18)	1.2034 (51)
H(16B)	0.7635 (43)	0.7254 (18)	1.0390 (54)
H(16C)	0.8269 (46)	0.6909 (17)	1.2172 (58)

motion parameters are listed in Table 3; no corrections for librational motion were applied to the positions or to the bond distances and angles given. The linear absorption coefficient is small (2.48 cm^{-1}), and absorption corrections were not applied.

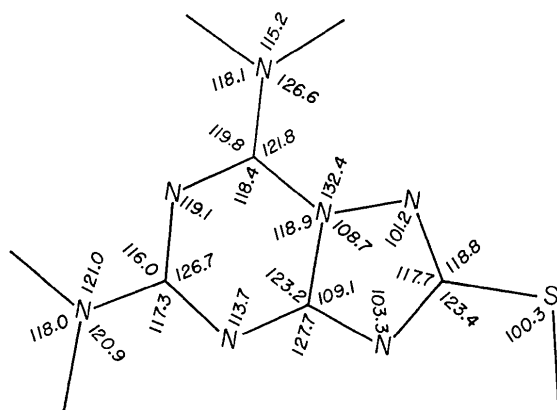


Fig. 2. Bond angles for (II). The standard deviations are 0.2° for angles involving the sulfur atom and 0.3° for all other angles involving non-hydrogen atoms.

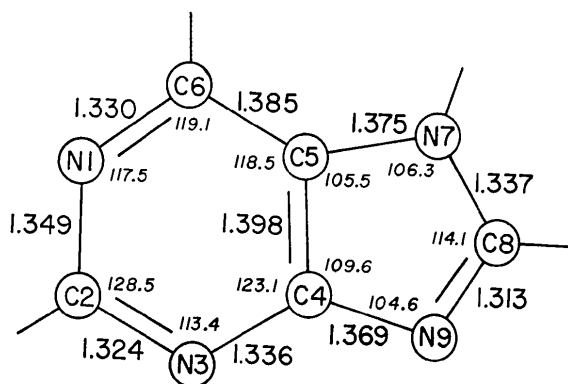


Fig. 3. Bond distances and angles of unsubstituted purine (Watson, Sweet & Marsh, 1965). Two sets of nearly identical parameters were presented by Watson, Sweet & Marsh; those of Sweet & Marsh are presented here.

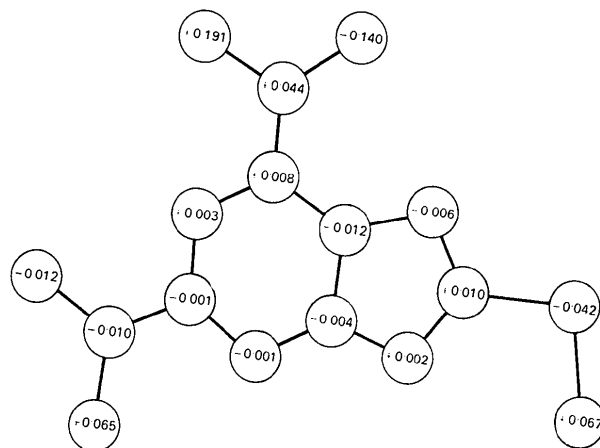


Fig. 4. Deviations of each atom in (II) from a best least-squares plane calculated through the atoms of the heterocyclic ring system.

Molecular structure

The triazolo-triazine ring system of (II) bears a strong resemblance to purine. Both are planar (to within approximately 0.01 \AA), and in both cases the observed bond distances indicate that the ten π electrons available on ring atoms in either system are extensively delocalized. The bond distances and angles of (II) are presented in Figs. 1 and 2. For comparison, the bond distances, bond angles, and formal bonding diagram of unsubstituted purine (Watson, Sweet & Marsh, 1965) are presented in Fig. 3. Although formally purine has a pyrimidine ring with three double bonds, and the triazolo-triazine derivative has an *s*-triazine ring with two double bonds, there are actually no significant differences in the parameters of the six-membered rings of the two compounds. In the five-membered rings of each, there are shortened bonds which indicate some degree of localization of double bond character between atoms which are formally double-bonded in the structural diagrams (II and Fig. 3). However, none of the ring bonds is as short as an isolated double bond or as long as an isolated single bond. The dimethyl-amino nitrogen atoms, N(10) and N(13), apparently contribute to the π -bond system also; they are connected to the ring system with bonds of hybrid length, and their bonding geometry is essentially planar.

The deviations of all atoms from a best least-squares plane calculated through just the ring atoms are shown in Fig. 4. Most of the distortions from planarity are minor, and could reasonably be attributed to intermolecular forces and/or fairly distant intramolecular repulsions. However, in the linked region, N(1)–N(8)–C(7)–N(10)–C(11), van der Waals repulsion forces should be large, for there are two consecutive *cis* junctions along this chain. There is a close approach between N(1) and C(11) of $2.893(5) \text{ \AA}$ compared to an estimated van der Waals approach of approximately 3.3 \AA . Also, one of the hydrogen atoms on C(11),

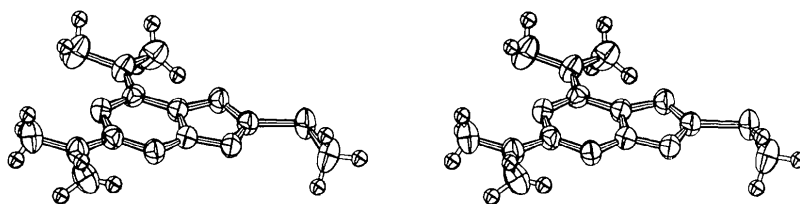


Fig. 5. A stereodrawing of (II), showing thermal vibration ellipsoids corresponding to the final thermal factors and isotropic hydrogen atoms at their refined positions.

H(11B), is situated only 2.27 (3) Å from N(1), compared to a van der Waals contact of 2.75 Å. The close approaches in this region would be considerably smaller if the N(10) dimethylamino group had exactly the same conformation as the N(13) dimethylamino group. The C(7)–N(10)–C(11) angle is 126.6 (3)°, more than 5° larger than the corresponding angles in the other dimethylamino group. Calculation of a plane through C(7), C(11), and C(12) reveals that N(10) is very slightly pyramidal; the N atom is 0.027 Å above the plane. The out-of-plane twist* of the N(10) dimethylamino group is 7.6° versus only 1.9° for the other one. Sundaralingam & Arora (1969) report that an out-of-plane twist of approximately 8° occurs in puromycin, which has a dimethylamino substituent in an analogous position on a purine ring system; full structural details were not published, so other structural adjustments in this vicinity cannot be compared.

Fig. 5 is a stereo-illustration which portrays the overall conformation of the molecule. The carbon, nitrogen and sulfur atoms are represented by thermal vibration ellipsoids calculated from the final thermal parameters, which are scaled to include 50% of the probability density.

Figs. 5, 6 and 7 were prepared with the help of the ORTEP computer program (Johnson, 1965).

Packing

The structure as viewed down **b** is shown in Fig. 6. The sheets of planar molecules are stacked parallel to the (10 $\bar{1}$) planes, with an interplanar distance of 3.36 Å. This stacking distance is similar to values observed in a wide variety of purines; e.g., in purine itself the distance is 3.39 Å (Watson, Sweet & Marsh, 1965). Two sheets of molecules are viewed down a direction which is perpendicular to their plane in Fig. 7, illustrating the overlap of the heterocyclic systems. Within each layer of this system, there is no opportunity for H-bond formation because each amino hydrogen has been replaced by a dimethylamino group. There is only one hydrogen-heavy atom approach that is shorter than a normal van der Waals contact. A methylthio hydrogen,

* The twist angle is defined as the average of two torsion angles; namely, $\omega(N6-C7-N10-C12)$ and $\omega(N8-C7-N10-C11)$. This is in accord with the twist angle defined by Winkler & Dunitz (1971) in their discussion of torsions in a semi-planar system.

H(16C), occurs only 2.58 (3) Å from N(4) of an adjacent coplanar molecule; the corresponding van der Waals approach is estimated to be about 2.74 Å. The shortest distance between stacked molecules is 3.360 Å, occurring between C(7) and N(4). There are, of course, several C–N and C–C contacts in the 3.4 to 3.5 Å range.

References

- CROMER, D. T. & LIBERMAN, D. (1970). *J. Chem. Phys.* **53**, 1891–1898.

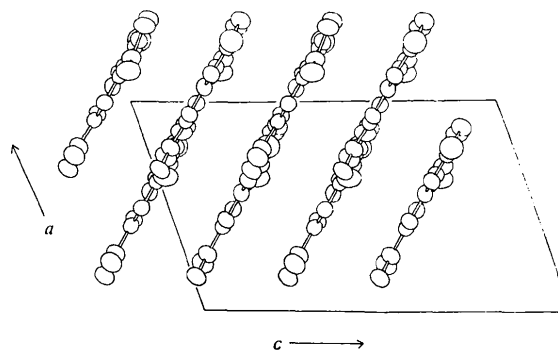


Fig. 6. A view of the packing down **b**. Hydrogen atoms are omitted for clarity; there are no hydrogen-hydrogen contacts between molecules less than 2.5 Å.

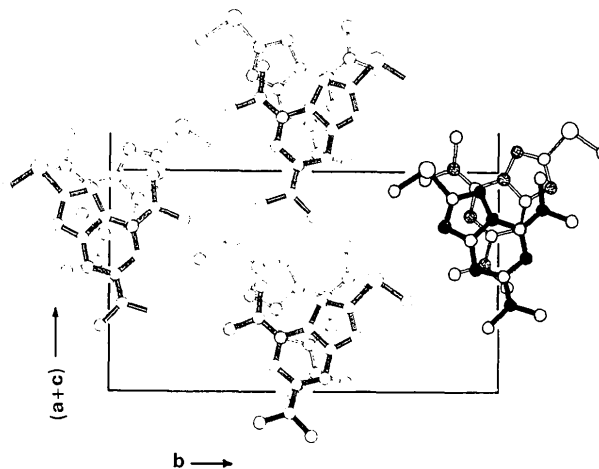


Fig. 7. A view of the packing normal to the layers shown in Fig. 6. Molecules from two layers are shown. In the pair of molecules at the right, nitrogen atoms are shaded and carbon atoms are left blank.

- DEMILO, A. B., OLIVER, J. E. & GILARDI, R. D. (1973). *J. Heterocyclic Chem.* **10**, 231-233.
- HAUPTMAN, H. & KARLE, J. (1959). *Acta Cryst.* **12**, 93-97.
- International Tables for X-ray Crystallography* (1962). Vol. III. Birmingham: Kynoch Press.
- JOHNSON, C. K. (1965). *ORTEP*. Report ORNL-3794, Oak Ridge National Laboratory, Oak Ridge, Tennessee.
- KARLE, J. (1968). *Acta Cryst.* **B24**, 182-186.
- KARLE, J. & HAUPTMAN, H. (1961). *Acta Cryst.* **14**, 217-223.
- KARLE, J. & KARLE, I. L. (1966). *Acta Cryst.* **21**, 849-859.
- LARSON, A. (1967). *Acta Cryst.* **23**, 664-665.
- SLETTEN, E., SLETTEN, J. & JENSEN, L. H. (1969). *Acta Cryst.* **B25**, 1330-1338.
- STEWART, R. F., DAVIDSON, E. R. & SIMPSON, W. T. (1965). *J. Chem. Phys.* **42**, 3175-3187.
- SUNDARALINGAM, M. & ARORA, S. K. (1969). *Proc. Natl. Acad. Sci.* **64**, 1021-1026.
- WATSON, D. G., SWEET, R. M. & MARSH, R. E. (1965). *Acta Cryst.* **19**, 573-580.
- WINKLER, F. K. & DUNITZ, J. D. (1971). *J. Mol. Biol.* **59**, 169-182.

Acta Cryst. (1973). **B29**, 2095

Structure Cristalline du Fluorosélénure d'Yttrium Orthorhombique YSeF 10

PAR NGUYEN HUY-DUNG

Laboratoire de Physique, Faculté des Sciences Pharmaceutiques et Biologiques de Paris-Luxembourg 4, Avenue de l'Observatoire, 75270 Paris Cedex 06, France

(Reçu le 14 mars 1973, accepté le 3 mai 1973)

The 10 form of yttrium fluoroselenide, YSeF, crystallizes in the orthorhombic system, with the space group *Pnam* ($a=9.912$, $b=6.314$, $c=4.080$ Å, $Z=4$). A single crystal was used for the structural determination by the heavy-atom method, and the atomic positions were refined by the least-squares method to $R=0.06$. The selenium atoms are at the vertex of a slightly deformed octahedron and have fourfold coordination while the three-coordinated fluorine atoms are at the centre of an yttrium isosceles triangle. The yttrium atoms have sevenfold coordination. The main feature in this structure is that there are plane groups of Y-F along two diagonals of the projected selenium hexagons. YSeF 10 is the basic structure of many polytype crystals in the Y-Se-F system.

Introduction

L'étude systématique des composés de lanthanides à deux anions a conduit récemment Dagron (1972) à isoler et décrire l'existence d'une famille de fluorosélénures d'yttrium et de lanthanides lourds, de formule générale LSeF (L=Ho, Er, Y). Les multiples phases observées dans le système Y-Se-F ont pu être identifiées sans ambiguïté, grâce à des études sur monocristaux. Ces composés sont tous des polytypes, dont la structure de base orthorhombique (que nous appelons YSeF10, d'après la notation de Ramsdell), était jusqu'à présent inconnue. Rappelons que Dagron a proposé un classement de ces polytypes en deux groupes. Le premier groupe de composés comprend, outre le monotype YSeF10, les polytypes orthorhombiques 30 et 70, dont les paramètres b sont des multiples 3 ou 7 du paramètre b_1 de la structure de base. Le deuxième groupe, en revanche, comprend les polytypes 2M, 3M, 4M, 5M, 6M, monocliniques. Leurs paramètres b vérifient la relation

$$b = nb_1 \sin \gamma \quad (n=2,3,4,5,6) .$$

Données expérimentales

Les cristaux de YSeF10 sont transparents, de couleur jaune pâle. Le cristal utilisé a une forme parallélépipédique de 80 μm de long sur 50 μm d'épaisseur et est allongé suivant l'axe c .

Les paramètres de la maille sont évalués à partir des clichés d'oscillation-rotation et de Weissenberg puis affinés par la méthode des moindres carrés à partir de quinze raies hkl d'un diffractogramme de Debye-Scherrer, avec un étalonnage donné par les raies du cuivre métallique: $a_1=9,912$ (4); $b_1=6,314$ (2); $c_1=4,080$ (1) Å. La densité expérimentale 4,86 g cm⁻³ conduit à placer quatre masses formulaires par maille, en accord avec la densité calculée (4,86 g cm⁻³).

Les réflexions obéissent aux conditions suivantes: $Ok1: k+l=2n$; $h0l: h=2n$. Les groupes spatiaux compatibles sont *Pnam* ou *Pna2₁*. Nous n'avons pas tenu compte de la présence de deux réflexions 902 et 052 d'intensité très faible dans la détermination du groupe spatial car leur existence observée dans le réseau réciproque d'un seul cristal est douteuse. Le calcul de la fonction de Patterson confirmera ultérieurement l'un de ces groupes spatiaux.